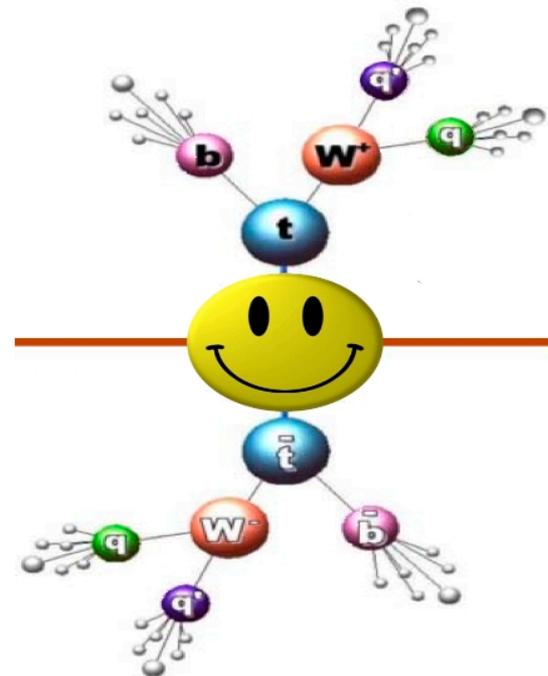
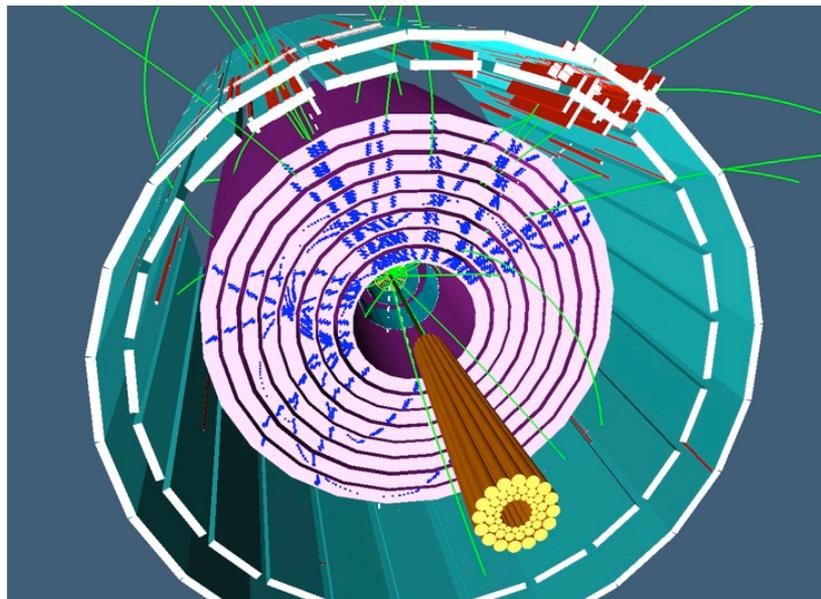




CLC

$$p\bar{p} \rightarrow X^0 \rightarrow t\bar{t}$$



Yuri Oksuzian  
University of Florida



# Outline I



- Service work
  - CLC Senior project leader and pager carrier
  - Testing and installation of the hardware components
  - HV CAEN crate maintenance
  - Studies on the PMT lifetime
  - PMTs gain instability
  - Development and implementation of single layer measurement
  - Monitoring and support of online luminosity measurement
  - Various tests and studies for joint luminosity meetings presentations
  - Ace and multiple CO shifts



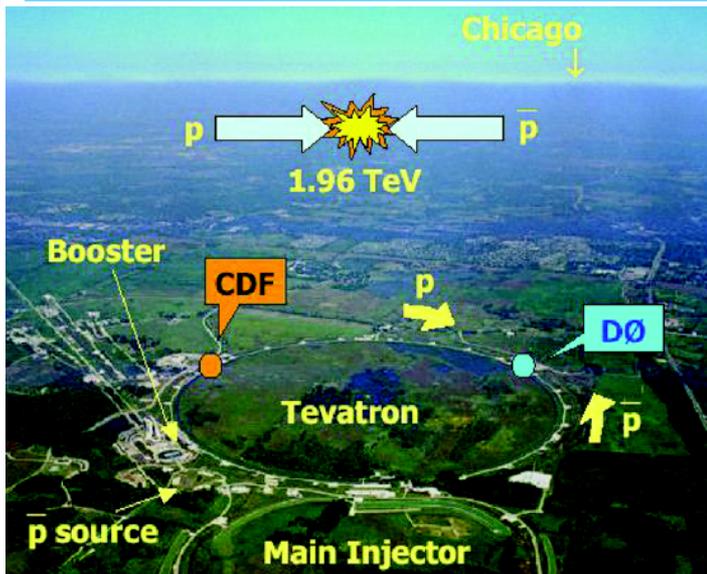
# Outline II



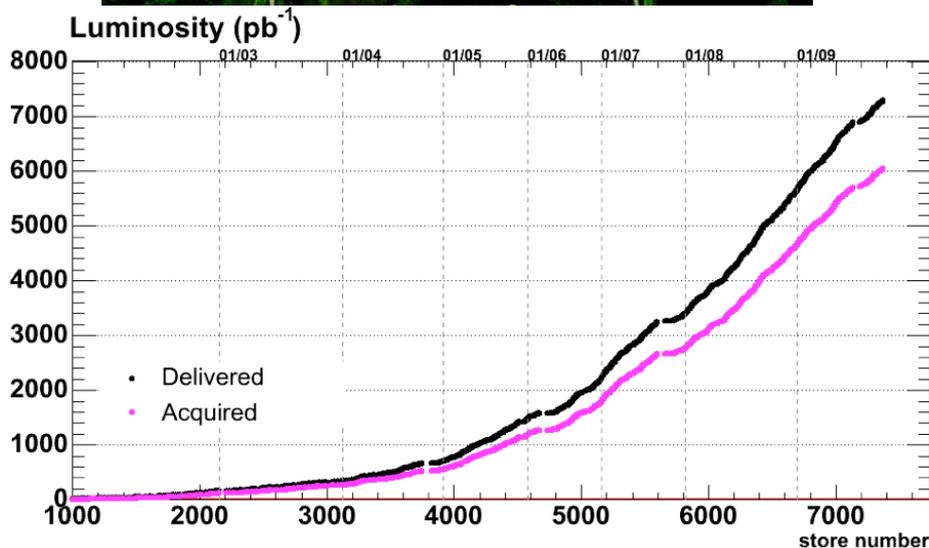
- Research
  - Search for non-standard model  $t\bar{t}$  resonance states in the all jets
  - FlaME implementation to all hadronic  $M_{t\bar{t}}$  search
  - Transfer functions optimization
  - QCD background modeling in all hadronic channel.  $M_{t\bar{t}}$  and  $M_{\text{mass}}$
  - Incorporation of events with extra jets
  - Neural net event selection optimization
  - FlaME implementation to event selection:  $M_{t\bar{t}}$  and  $M_{\text{mass}}$
  - Studies on combinatorial reweighing, using FlaME, JetProb.



# Tevatron



- Operating since 1985
- Highest energy collider
- Collides  $p\bar{p}$
- 36x36 bunches
- Average xring rate  $\approx 1.7\text{MHz}$
- Radius  $\approx 1\text{km}$
- $L_{\text{ints}} \approx 3.5 \cdot 10^{32} \text{cm}^{-2} \text{s}^{-1}$
- Particles are accelerated in 5 steps
- 7  $\text{fb}^{-1}$  already collected.
- 12  $\text{fb}^{-1}$ , by the end of Run II

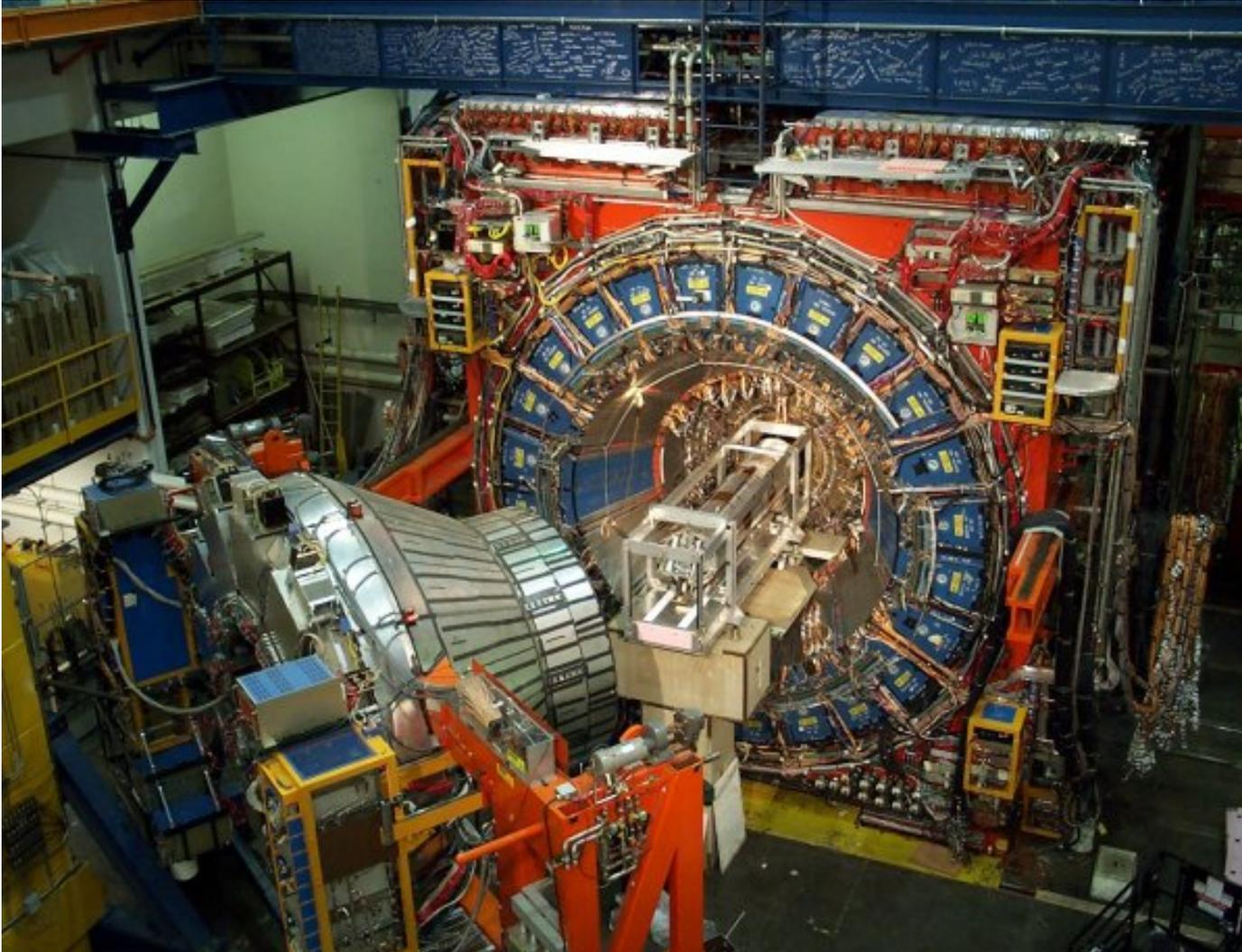


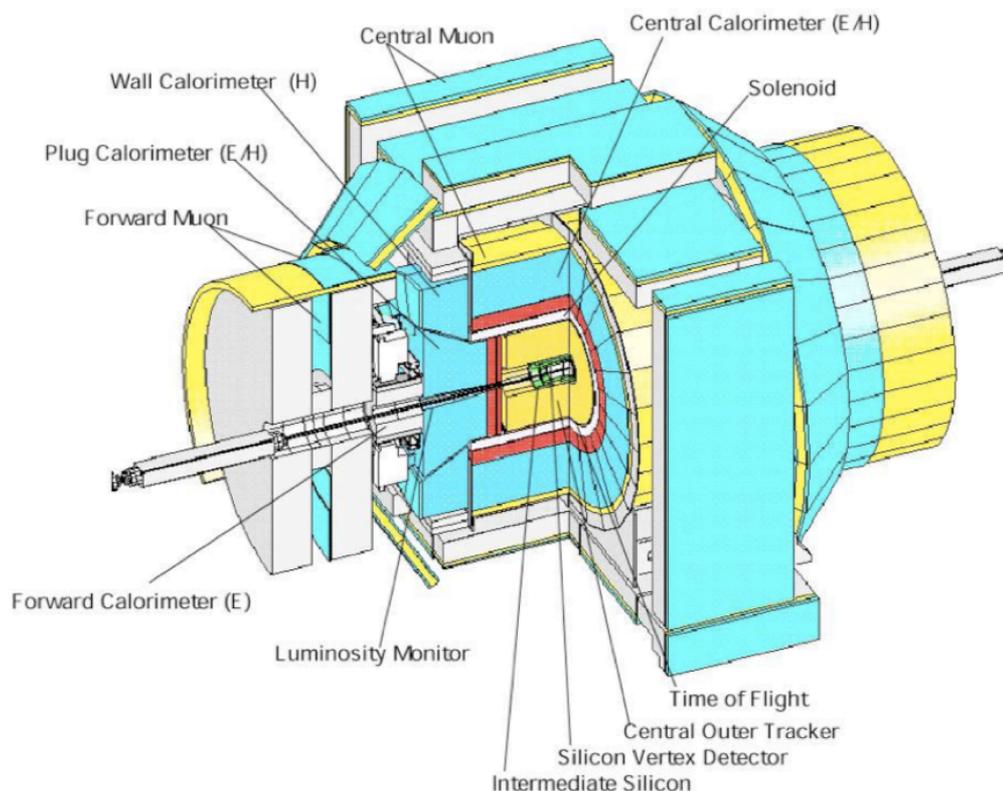


# CDF



10 meters high, 27 meters long and weighs over 5000 ton





- General purpose detector
- Need the full detector for my analyses:
  - Tracking
  - Calorimeter system
  - Muon system
  - Luminosity counter
- Tracking:
  - Silicon
    - L00+SVX+ISL
    - Precise vertex tracking
    - $R \approx 1.6-28\text{cm}$
    - $|\eta| < 2.0/4.0$
  - COT
    - Drift chamber with 8 layers
    - Tracking in central region
    - $R \approx 44-132\text{cm}$
    - $|\eta| < 1.0$
- Calorimeter system
  - Central, wall and plug
  - Electron/photon/jet/ $\Sigma E_T$  energies
  - Overall  $|\eta| < 3.6$



# CLC

Cerenkov Luminosity Counters



# Zero crossing method



$$\mathcal{L} = \frac{f_{BC}}{\epsilon_{in}^{CLC} \cdot \sigma_{in}} \quad \mu$$

We need to calculate # of interactions per bunch crossing

Zero bunch crossing probability approximately

$$\mathcal{P}_0(\mu) = e^{-\mu}$$

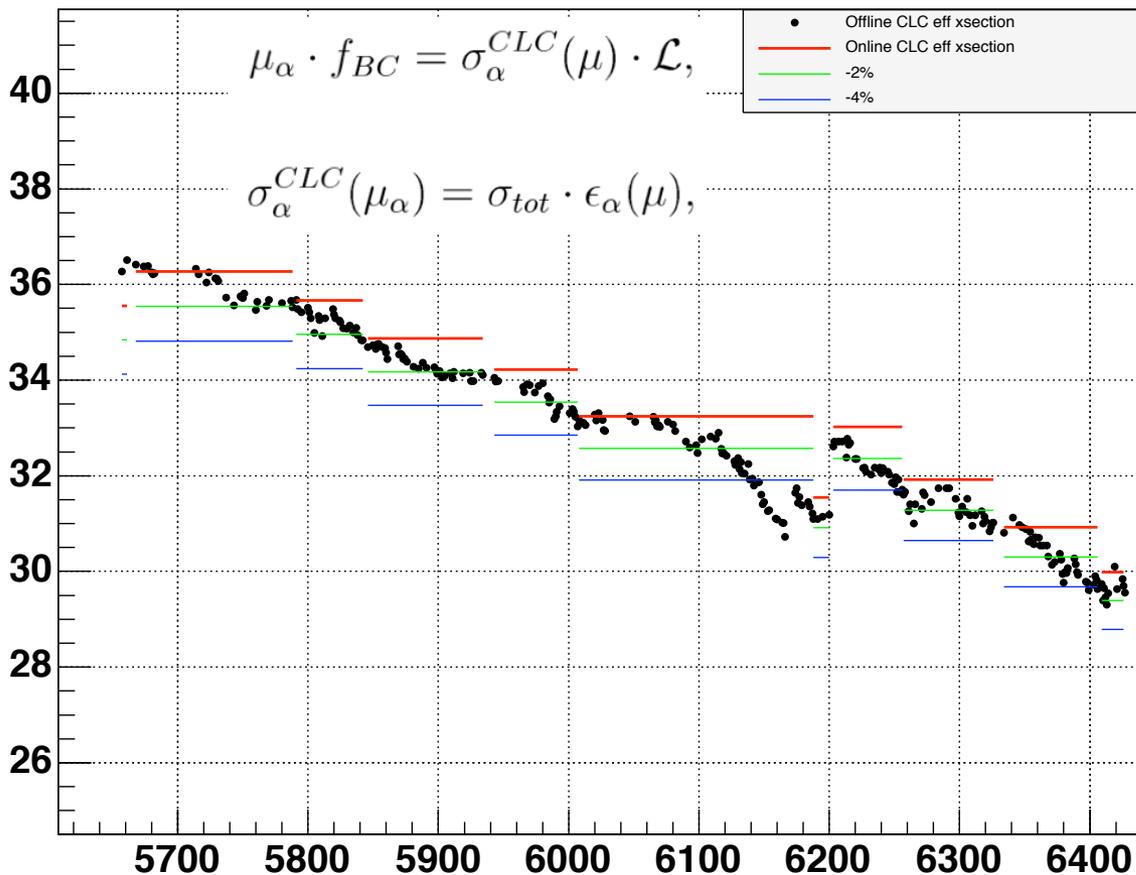
In reality, N # of interactions will not be detected independently

$$P_o^{exp}\{\mu; \alpha\} = \sum_{k=0}^{\infty} \mathcal{P}_k(\mu) \cdot \pi_0(k) = (e^{\epsilon_w \cdot \mu} + e^{\epsilon_e \cdot \mu} - 1) \cdot e^{-(1-\epsilon_0) \cdot \mu}$$

Trough CDF simulation, we calculate all the acceptances above, using CLC response, i.e. SPP values for each channel

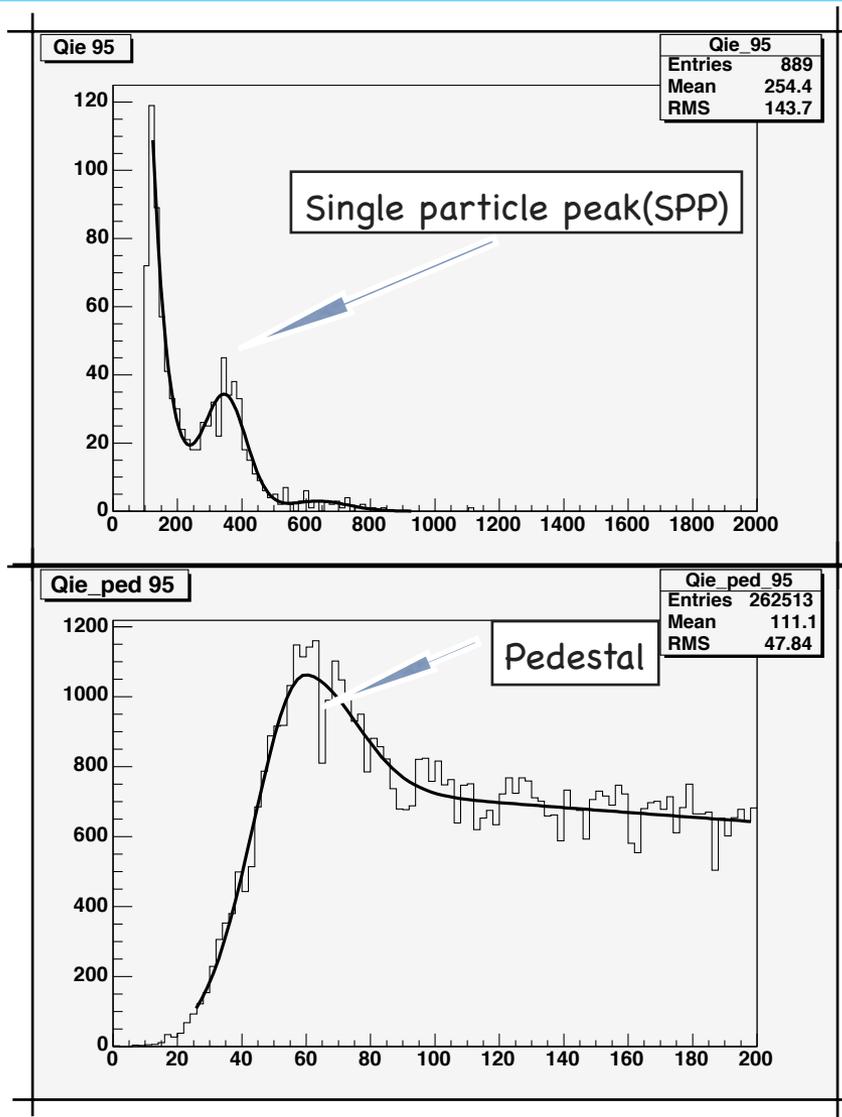


# Effective xsection changes



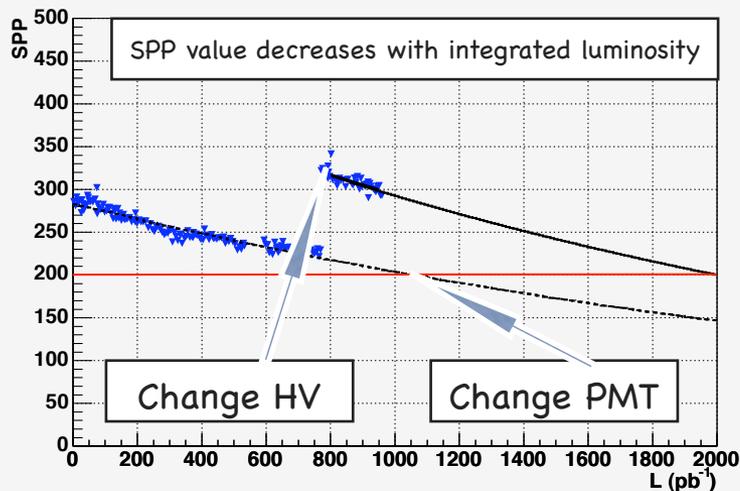


# PMT: SPP, Pedestal

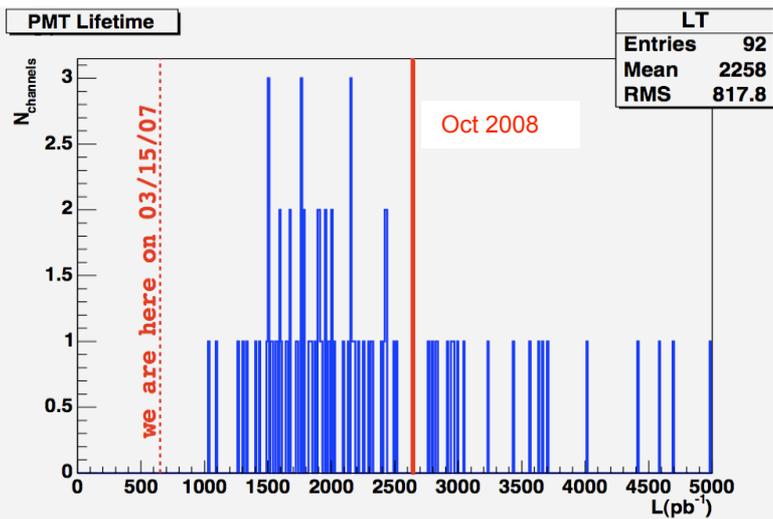




# PMT lifetime



Money saved

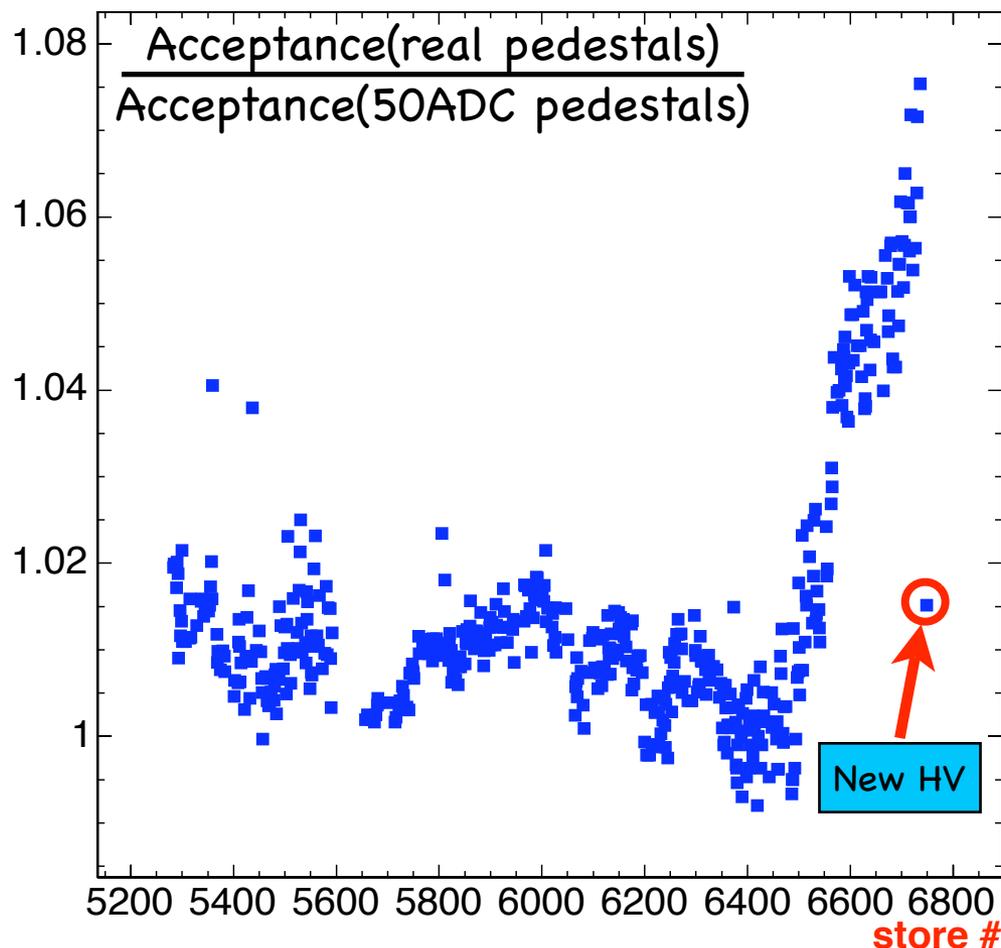




# Lower gain period



- From stores ~6500 (mid Oct 2008), we start gradually underestimate acceptance
- It directly translates into luminosity overestimation
- Last store with new HV settings shows only 1.5% difference



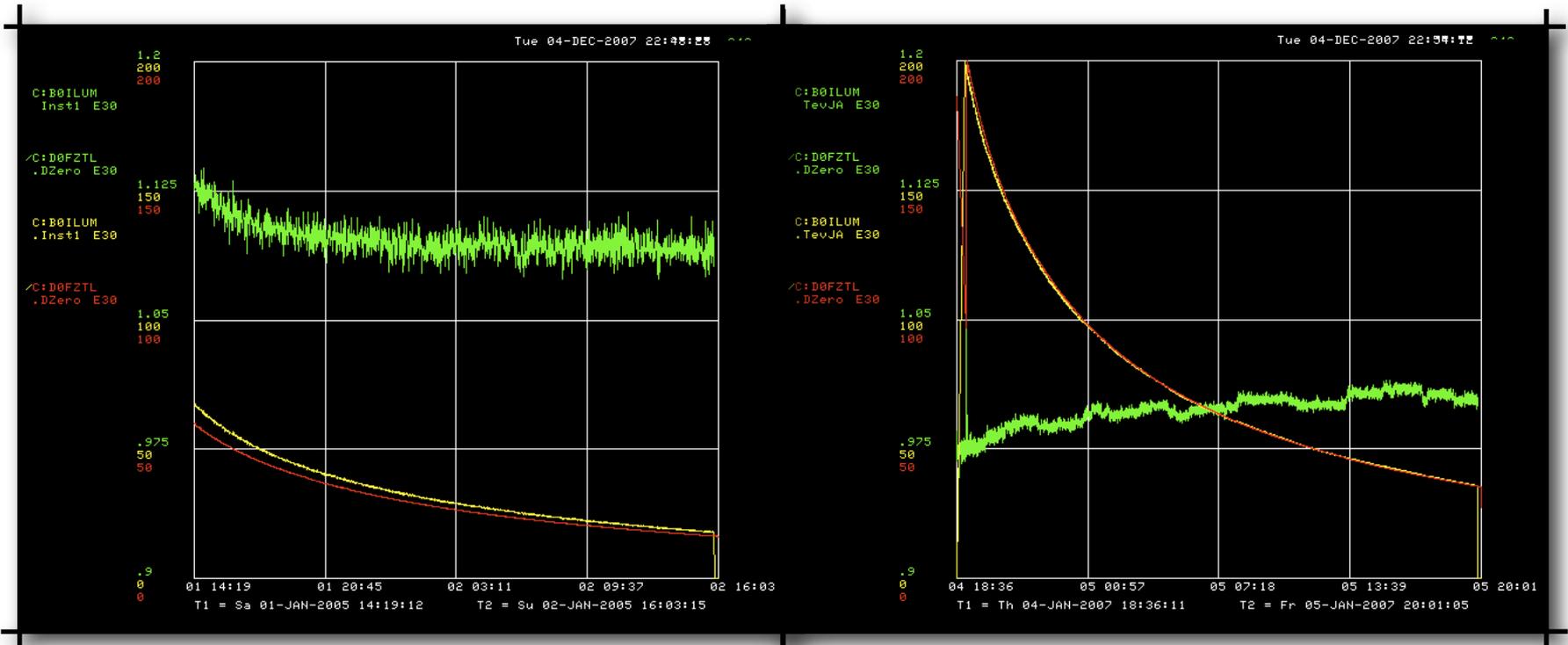


# CDF/D0



## 2005

## 2007





# Amplitude vs time

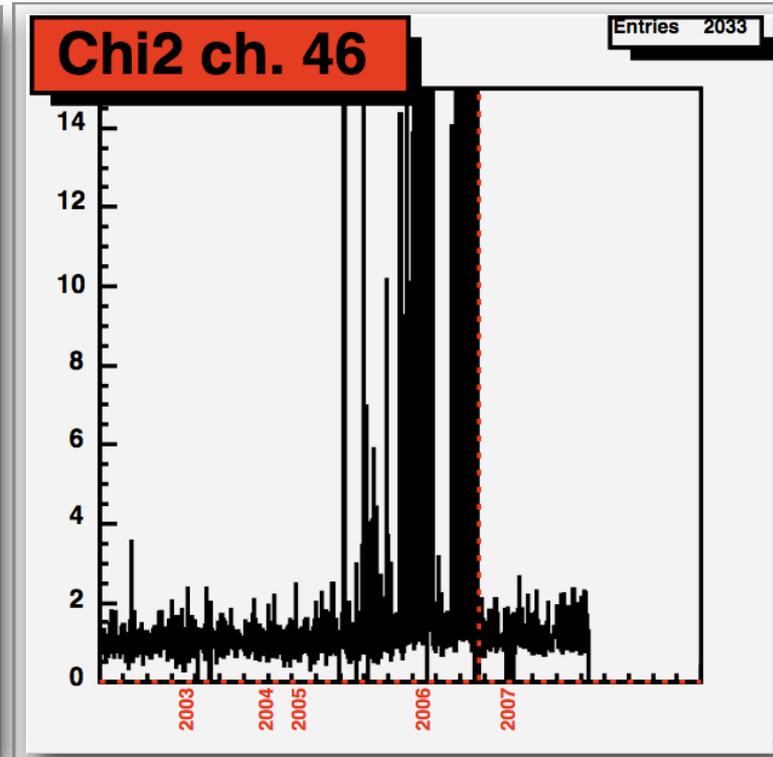
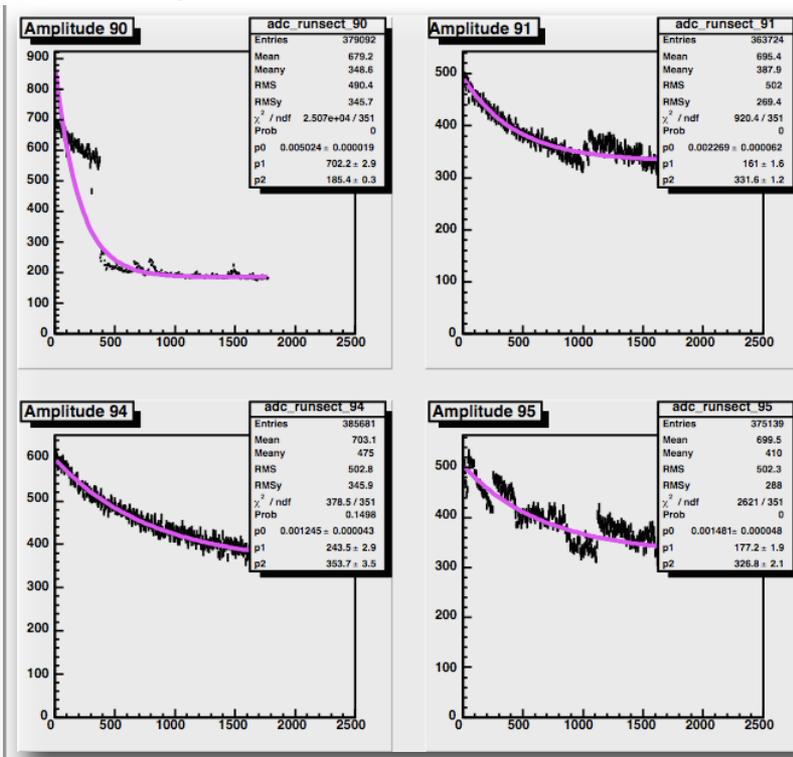


Here we plot:

- Y-axis: <ADC>
  - X-axis: Run Section
- <ADC> expected to fall smoothly during data taking

Here we plot:

- Y-axis: time
  - X-axis: Chi2 of the fit
- Chi2 ~1 for a stable channel



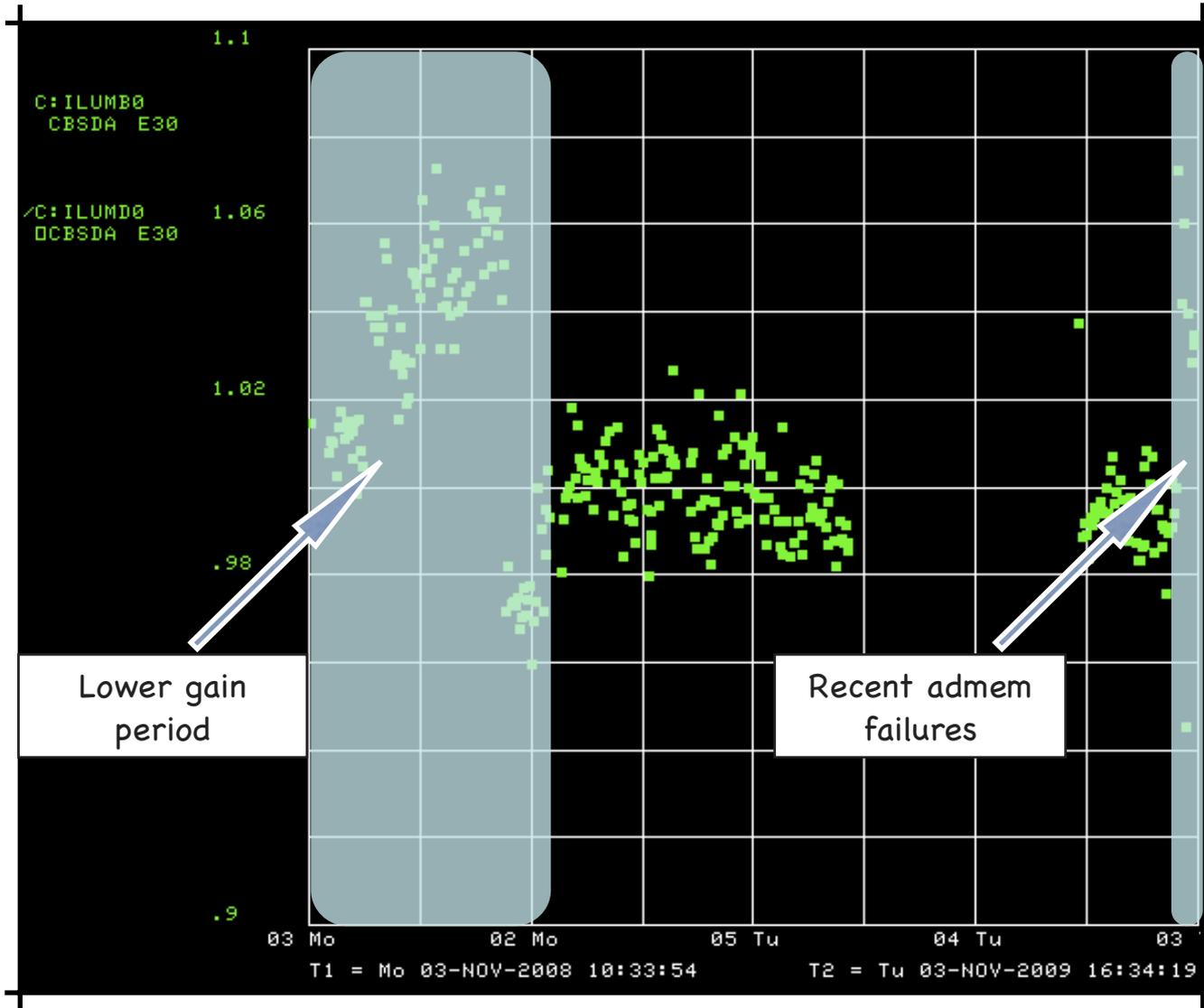
# Defective base



- Traced down the instability in luminosity measurement to bad soldering in PMT bases
- Replaced 44(~70%) bases.



# CDF/D0





# Analysis

$$p\bar{p} \rightarrow X^0 \rightarrow t\bar{t}$$



# Why and How?



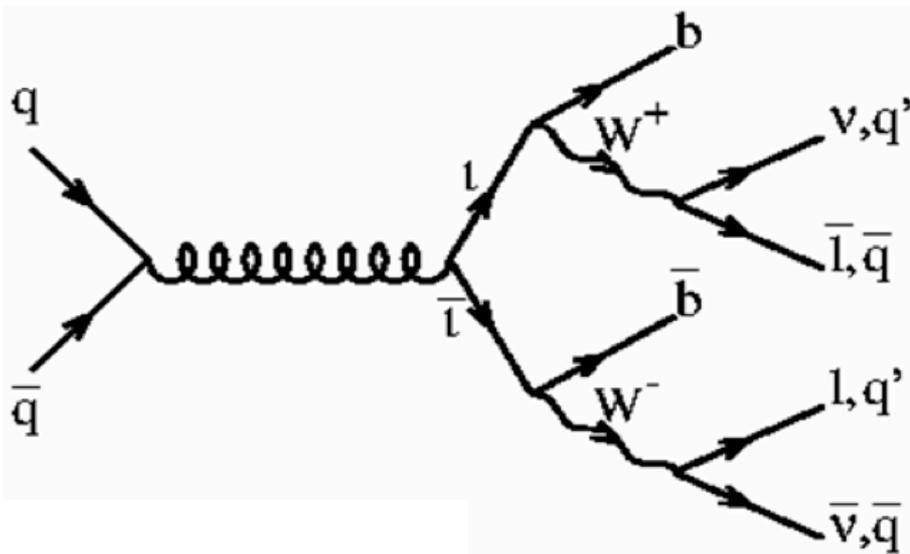
- Goal is to test  $t\bar{t}$  production for possible new sources such as a narrow resonance
  - Top is very heavy, maybe indication of coupling to new physics
  - Top is a young particle
  - Various theoretical models predict it
- Search technique:
  - $M_{t\bar{t}}$  spectrum is reconstructed, using FlaME
  - Search for a peak in  $M_{t\bar{t}}$  spectrum
    - Understand SM fluctuation probabilities
    - Calculate UL(Upper Limits)
    - Compare data with our expectations(SM or with new physics)



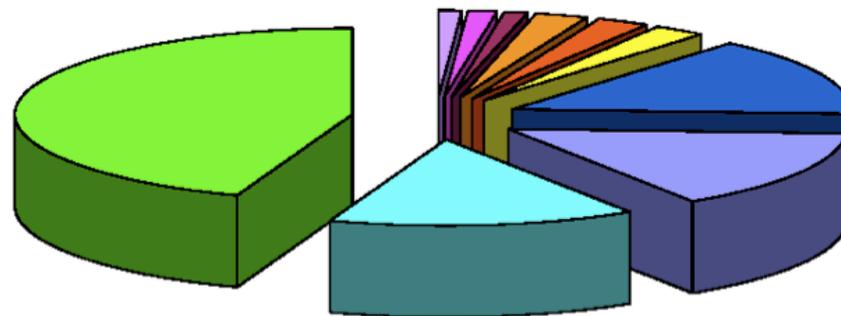
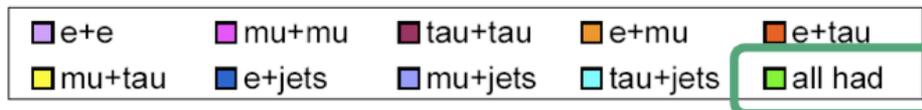
# Where?



- First  $M_{t\bar{t}}$  analyses in All Hadronic channel
  - Disadvantages
    - Large QCD background
      - » Controlled with good event selection
    - More combinations
  - Advantages
    - Highest branching ratio
      - » Most  $t\bar{t}$  events are here
    - No missing information like neutrino
      - » Better signal templates
  - Future
    - Combined result with lepton+jets channel
      - » Higher sensitivity
    - Cross-check for a possible discovery



ttbar Decay Modes



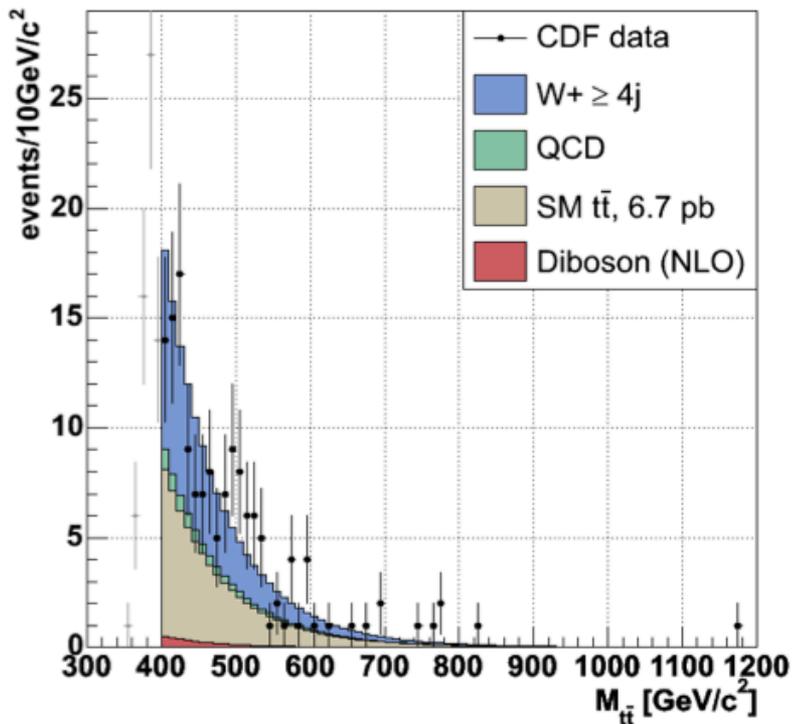
- Dilepton:
  - BR = 5% ( $ee/e\mu/\mu\mu$ ) - cleanest sample, lowest statistics
- Lepton+Jets:
  - BR = 29% ( $e/\mu$ +jets only) - golden channel with high statistics and reasonable S/B
- All-hadronic:
  - BR = 46% - highest statistics but large background



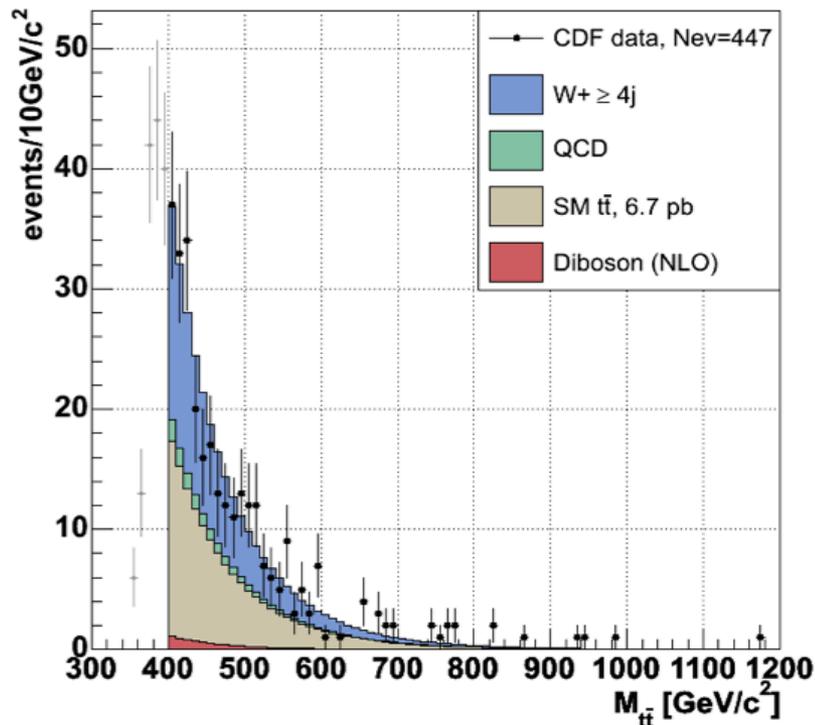
# Previous L+jets result



CDF Run 2 preliminary, L=319pb<sup>-1</sup>

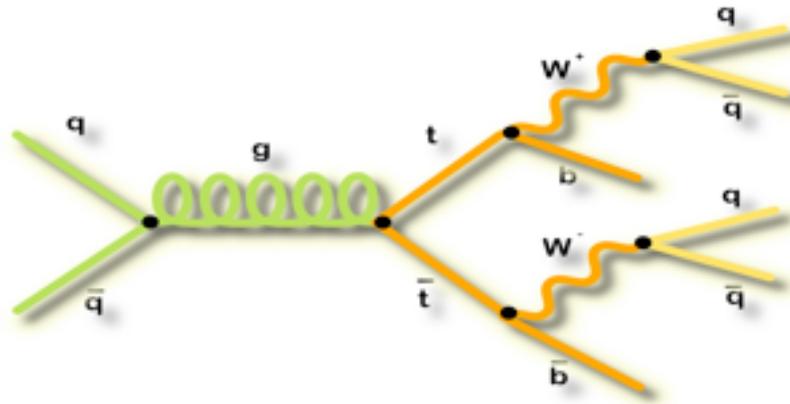


CDF Run 2 preliminary, L=682pb<sup>-1</sup>





# FlaME (Florida Matrix Element)



We calculate the *a priori* probability density for an event to be the result of Standard Model  $t\bar{t}$  production and decay

$$P(j | M_{top}) = \frac{1}{\sigma(M_{top})\epsilon(M_{top})N_{combi}} \sum_{combi} \int dz_a dz_b f(z_a) f(z_b) d\sigma(M_{top}, p) TF(j | p) P_T(p)$$

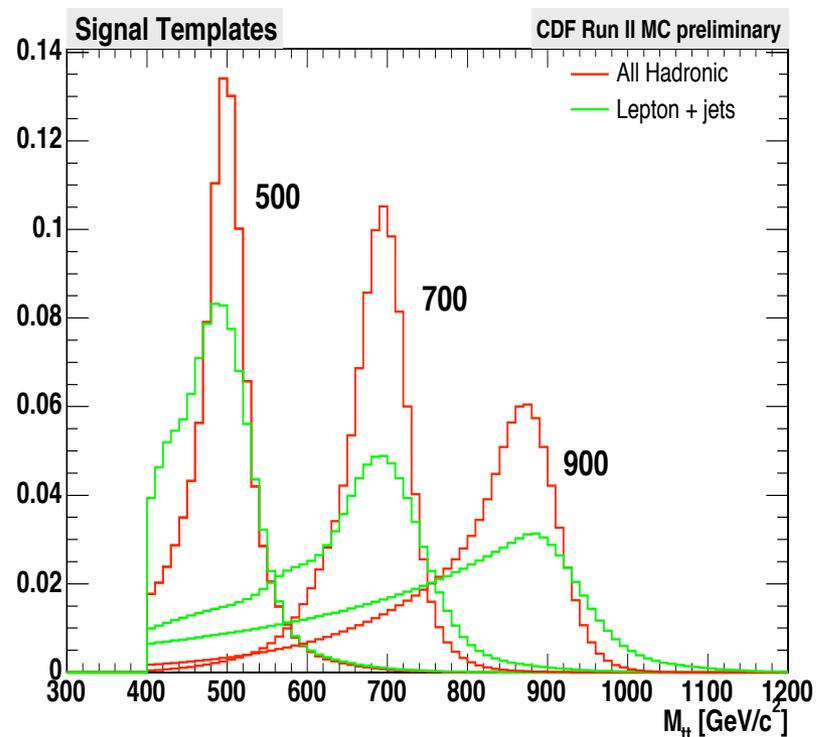
To calculate the  $M_{t\bar{t}}$  probability density, we modify the integral above:

$$\rho(x | j) = \frac{1}{\sigma(M_{top})\epsilon(M_{top})N_{combi}} \sum_{combi} \int dz_a dz_b f(z_a) f(z_b) d\sigma(M_{top}, p) TF(j | p) P_T(p) \delta(x - M_{t\bar{t}}(p))$$

As  $M_{t\bar{t}}$  estimator we use average of this distribution:

$$M_{t\bar{t}} = \langle \rho(x | j) \rangle$$

- Signal samples:
  - Pythia generated narrow resonant  $t\bar{t}$  samples with masses 450, 500 ... 900 GeV
- Background Samples:
  - SM  $t\bar{t}$  MC sample
  - QCD
    - Data driven





# Trigger & Prerequisites



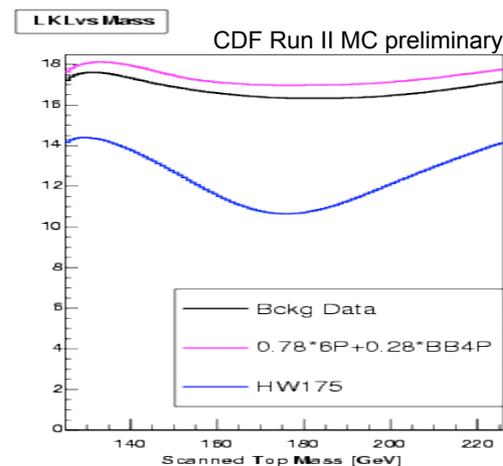
- Multijet Trigger
  - L1:  $\geq 1$  tower with  $E_T \geq 10$  GeV
  - L2:  $\geq 4$  clusters with  $E_T^{cl} \geq 15$  GeV,  $\Sigma E_T \geq 125$  GeV
  - L3:  $N_{jet} \geq 4$ , with  $E_T^{jet} \geq 10$  GeV
    - $\sigma \approx 14$  nb,  $\sim 85\%$  all hadronic efficiency
- Prerequisites
  - Good run list
  - Vertex:  $|z| < 60$ cm &  $|z - z_p| < 5$ cm
  - Missing Et Significance:  $< 3$  (GeV)<sup>1/2</sup>
  - Tight lepton veto
  - 6,7 tight jets -  $E_T^{jet} \geq 15$ GeV,  $|\eta| < 2.0$
- After prerequisites we have  $t\bar{t}/\text{QCD} \sim 1/1000!$



# Neural Net Idea



- Neural net event selection:
  - Uses a Root class TMultiLayerPerceptron
  - 11 inputs, 2 hidden layers with 20/10 nodes and 1 output
- SumEt - total transverse energy
- SumEt3 - sub-leading transverse energy  $\sum E_T - E_{T1} - E_{T2}$
- C - centrality:  $\sum E_T / \sqrt{\hat{s}}$
- A - aplanarity:  $3/2 * (\text{smallest eigenvalue})$  of  $M_{ab} = \sum_j P_a^j P_b^j / \sum_j |\vec{P}^j|$
- $E_N^*$  - geom average of transverse energy of the N-(2 leading jets)
- $E_{T1}^*$  - transverse energy of the leading jet
- $M_{2j}^{\min}$  - the minimum dijet mass
- $M_{2j}^{\max}$  - the maximum dijet mass
- $M_{3j}^{\min}$  - the minimum trijet mass
- $M_{3j}^{\max}$  - the maximum trijet mass
- FlaME variable,  $\sum -\text{Log}(P(M_{\text{top}}=155, 160 \dots 195 \text{ GeV}))$





# QCD background



- We build tag matrix from events from 4,5 jet events.
- Each element in the matrix defined as:

$$P^+(N_{v12}, E_{tjet}, N_{tr,jet}) \equiv \frac{N_{j,tag}(N_{v12}, E_{tjet}, N_{tr,jet})}{N_{j,tgbl}(N_{v12}, E_{tjet}, N_{tr,jet})}$$

- The probability to single/double tag an event:

$$\sum_i \left[ P_i^+ \cdot \prod_{j \neq i} (1 - P_j^+) \right] \quad \sum_{i,j \neq i} \left[ P_i^+ \cdot P_j^+ \cdot \prod_{k \neq i,j} (1 - P_k^+) \right]$$

- We weigh each event in pre-tagged data sample to get the prediction for 1, 2 tagged events
- Finally, we define several control region and test our modeling with observation
- For all control regions we get a very good agreement
- Biggest impact on final result comes from possible signal contamination, using this procedure

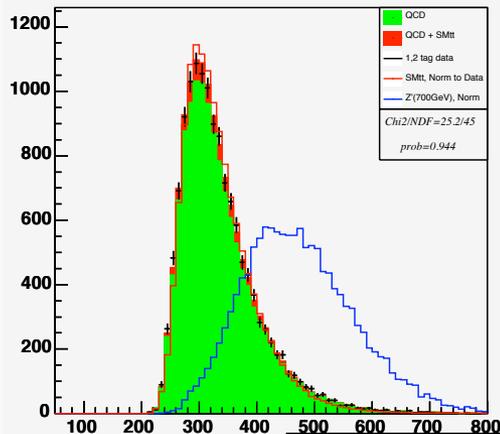


# Crosscheck with data

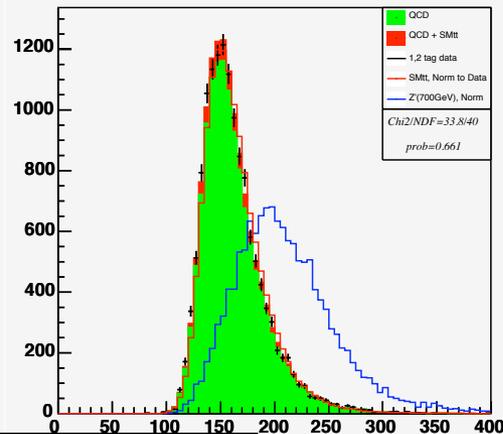


$0.75 < \text{NNetOut} < 0.93$ .

SumEt

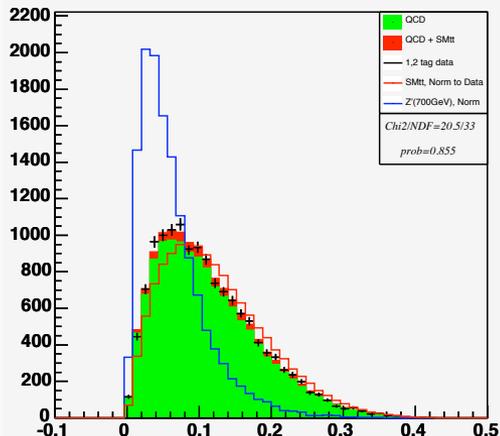


SumEt3

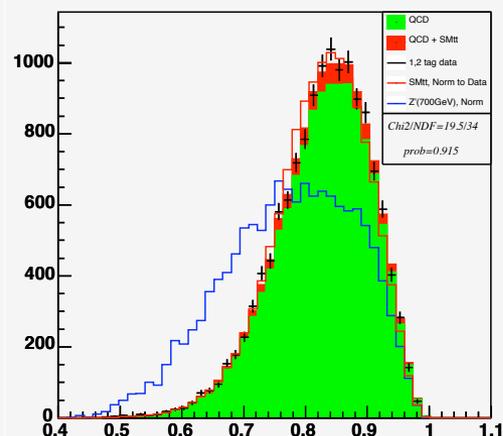


CDF Run 2 preliminary

Apla



Cent

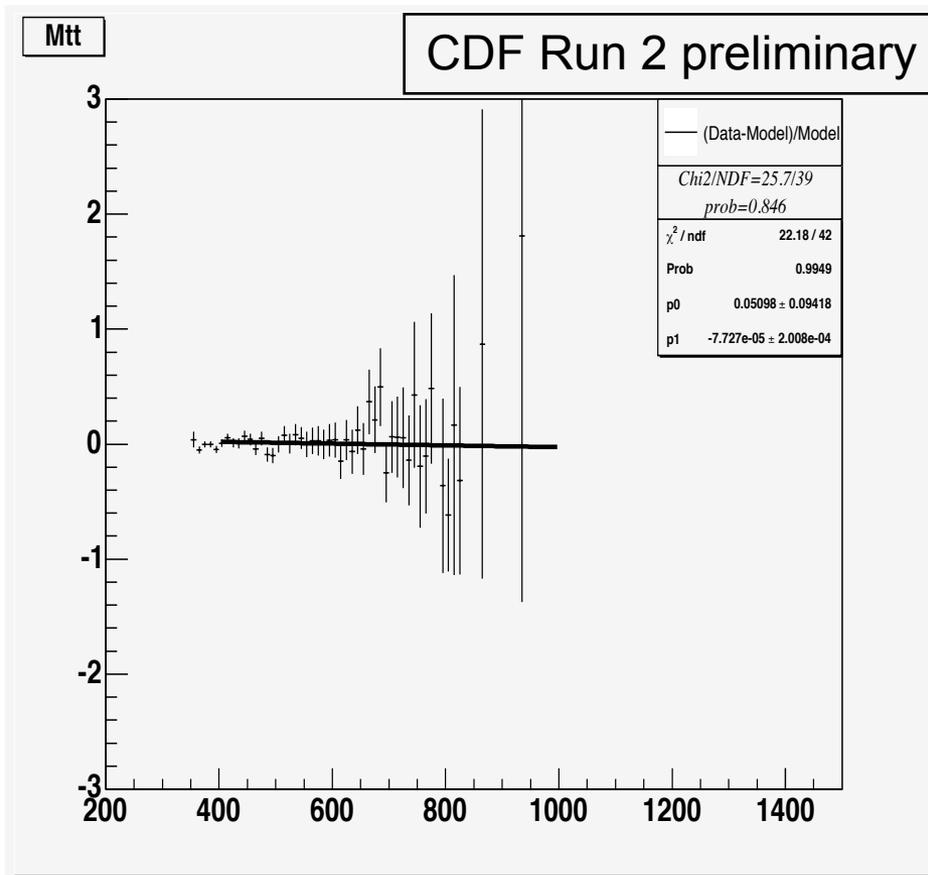
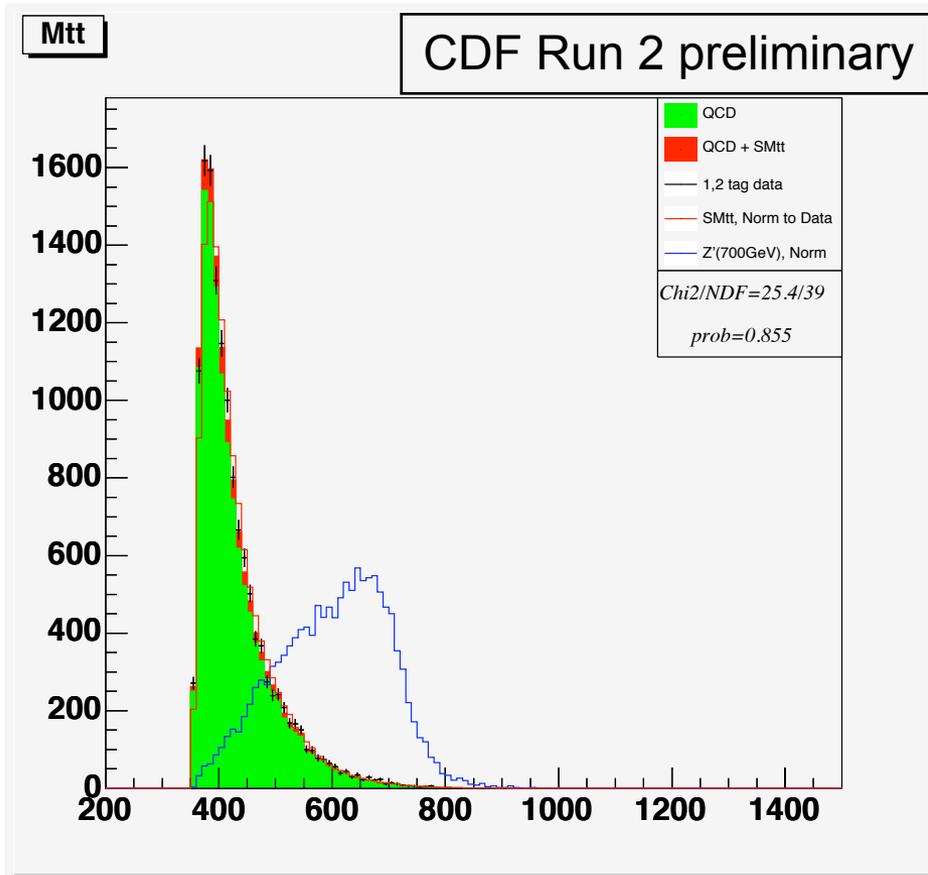




# Crosscheck with data



$0.75 < \text{NNetOut} < 0.93$ .





# Limit Setting Methodology



- Template event weighting
  - $N_{X0}$ : based on assumed cross-section and acceptance
  - $N_{tt}$ : based on theoretical cross-section and acceptance
  - $N_{QCD}$ : Balance from data

$$N_{cdf}^{tot} = \int L dt \cdot (\sigma_{X0} A_{X0} + \sigma_{t\bar{t}} A_{t\bar{t}}) + N_{QCD}$$

- Likelihood

- $N_{X0}$ ,  $N_{tt}$ ,  $N_{QCD}$  are used to compute the expected number of events in mass bin “i”:

$$\mu(i) = N_{X0} T_{X0}(i) + N_{tt} T_{tt}(i) + N_{QCD} T_{QCD}(i)$$

- Given the observed number of events  $n(i)$  and expected  $\mu(i)$  in bin “i”, the likelihood is equal to:

$$L(\sigma_{X0}, \vec{v} | \vec{n}) = \prod e^{-\mu_i} \frac{\mu_i^{n_i}}{n_i!}$$



# Posterior density function



- Acceptance uncertainties accounting

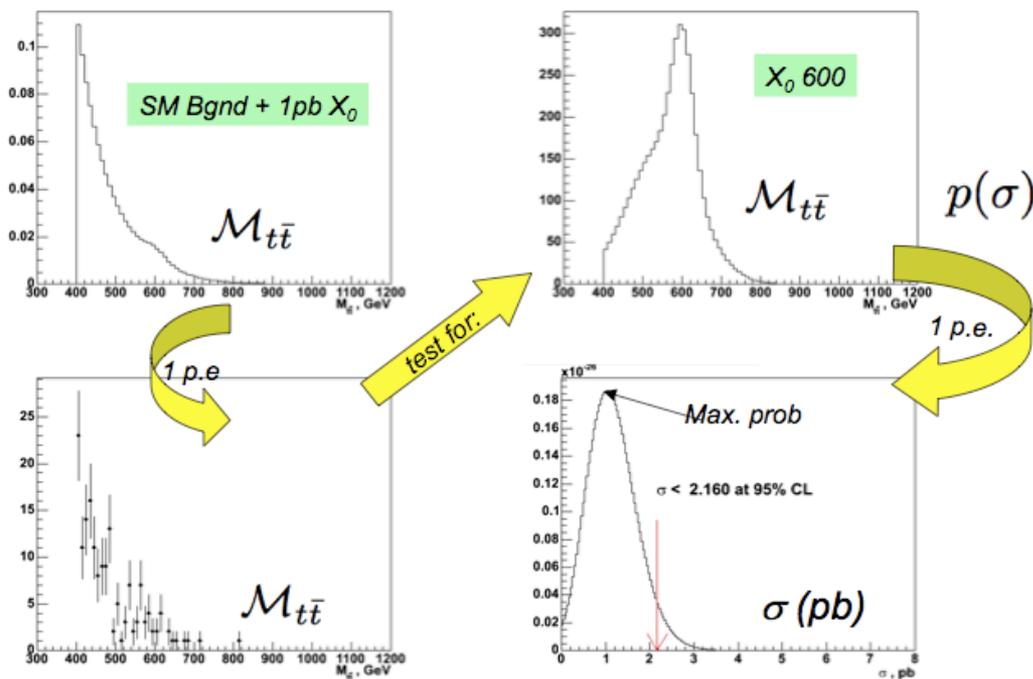
$$p(\sigma_{X_0}, \vec{n}) = \int d\vec{v} \cdot L(\sigma_{X_0}, \vec{v} | \vec{n}) \cdot \pi(\sigma, \vec{v})$$

- We integrate over the nuisance parameters, uncertainties for:

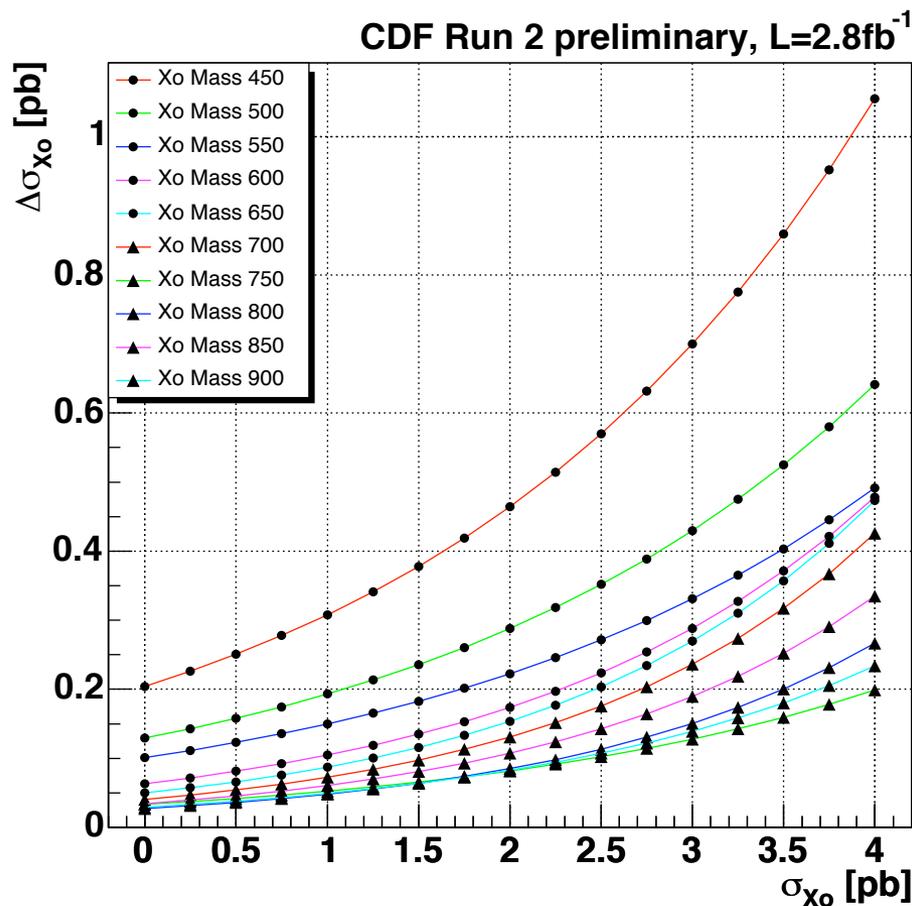
- Signal acceptance
- Background acceptance
- Background cross-section

- Given  $p(\sigma|\vec{n})$  we define:

- $\sigma_{X_0}$  - max of PDF
- 95% confidence level upper limit (UL)  $\frac{1}{Area} \int_0^{UL} p(\sigma | \vec{n}) d\sigma = 0.95$
- Values are calculated as median after 1000 PE's



- To consider systematics, which both affect shape and acceptances, we:
- Consider the shift on cross-section by:
- Running PE from shifted templates and fit them with nominal ones
- We considered systematics due to JES, ISR/FSR. PDF found to be negligible

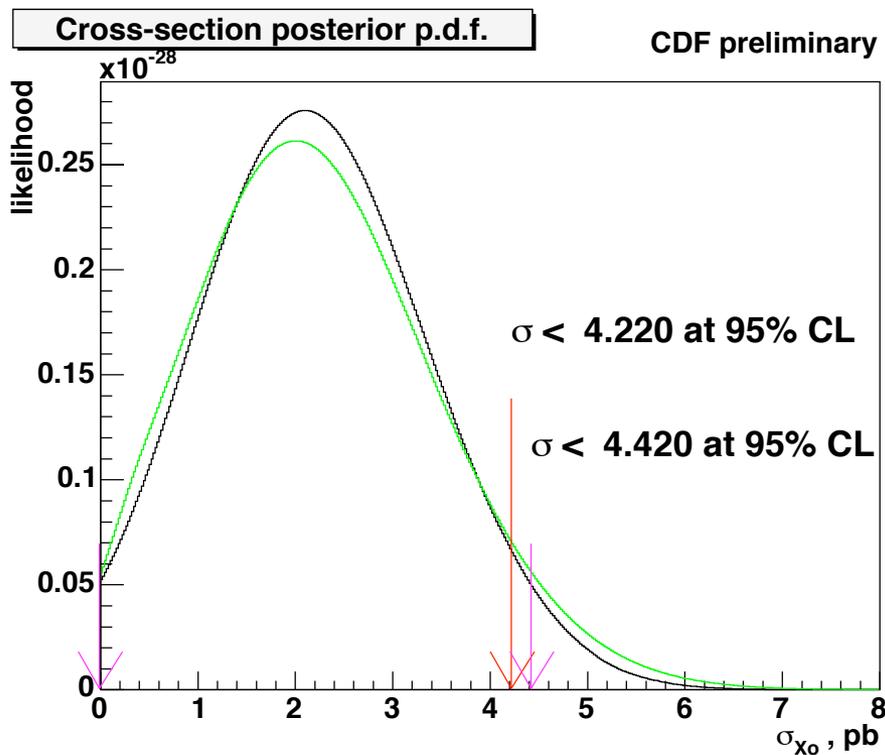




# Applying systematics

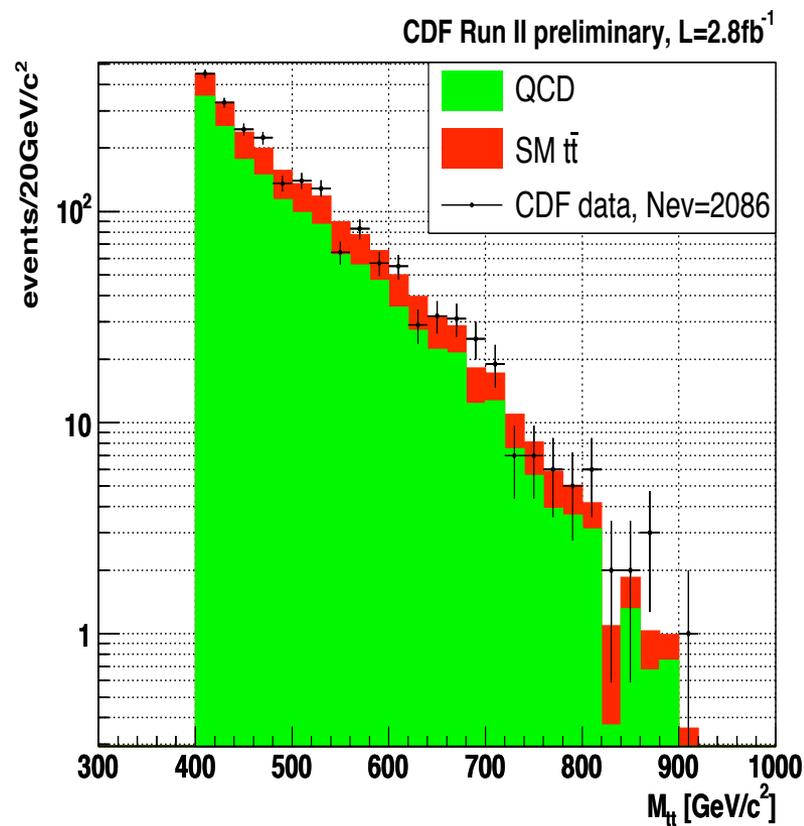
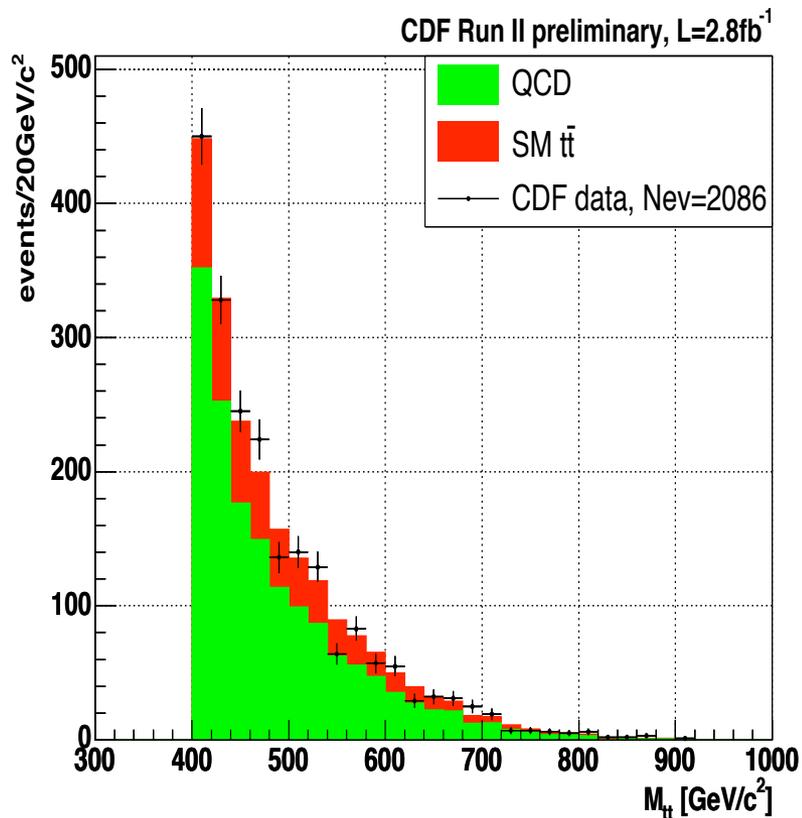


$$PDF_{SYS}(\sigma_{X_0}) = \int_0^{\infty} \frac{1}{\delta\sigma_{X_0} \sqrt{2\pi}} \exp\left(-\frac{1}{2} \left(\frac{\sigma_{X_0} - \sigma'}{\delta\sigma_{X_0}}\right)^2\right) PDF(\sigma') \cdot d\sigma'$$



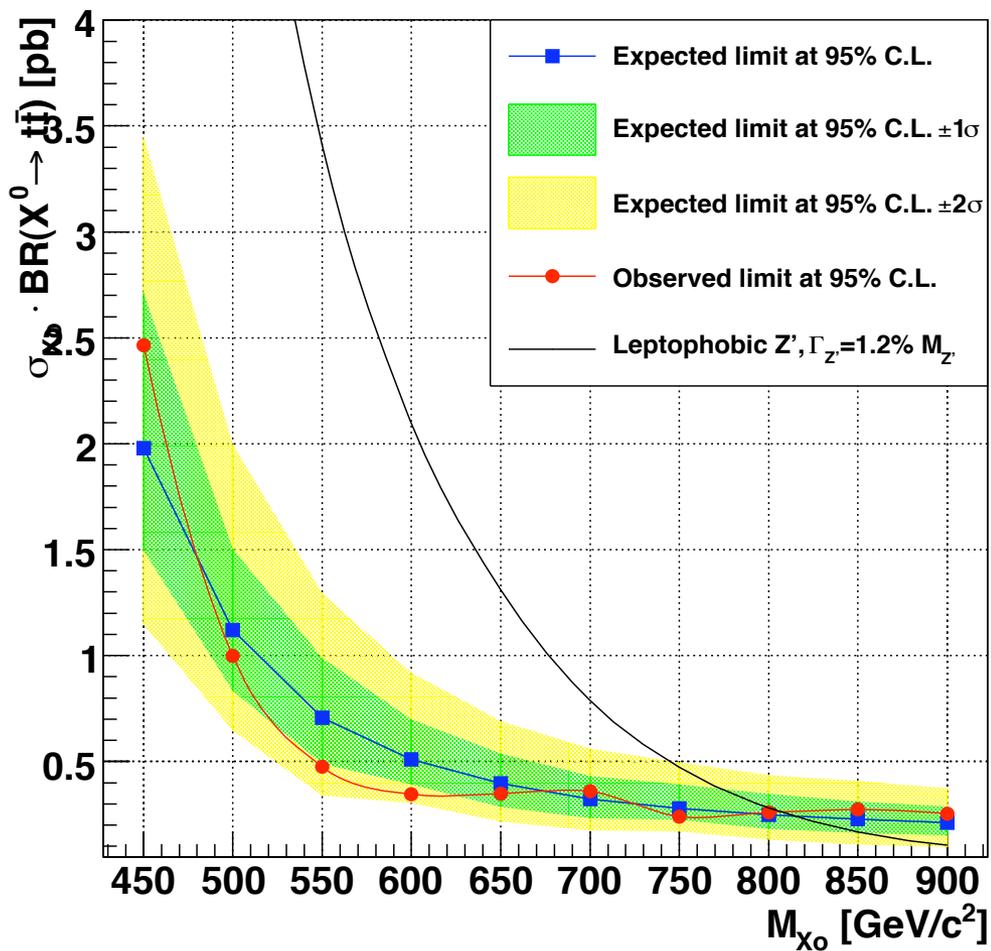


# Data/BG prediction





# Upper Limits





# Conclusions & Plans



- Maintenance and support for CLC
- First search for  $t\bar{t}b\bar{b}$  resonance in all jets final state
  - No excess found in 2.8/fb of CDF data
  - We set observed upper limit on leptophobic  $Z'$  mass up to 805 GeV
- Plans:
  - Graduate in December 2009
  - Join the strong group with
    - » Hardware projects
    - » Challenging topics in search analysis



# Backup slides



# Signal contamination



- Signal contribution to QCD shape will be treated as following:

From Equation 4 we have, number of events in bin “i”:

$$\mu = \sigma_s A_s T_s + \sigma_{tt} A_{tt} T_{tt} + N_{QCD}^{pure} T_{QCD}^{pure}$$

$$N_{QCD}^{pure} T_{QCD}^{pure} = N_{QCD}^{cont} T_{QCD}^{cont} - \sigma_s A_s^{cont} T_s^{cont} - \sigma_{tt} A_{tt}^{cont} T_{tt}^{cont}$$

Comparing signal templates of predicted and observed values we can assume:

$$T_s = T_s^{cont}$$

So, finally we get:

$$\mu = \sigma_s (A_s - A_s^{cont}) T_s + \sigma_{tt} A_{tt} T_{tt} + N_{QCD}^{cont} T_{QCD}^{cont} - \sigma_{tt} A_{tt}^{cont} T_{tt}^{cont}$$

- In the end, it decreases signal acceptances by the values we get from TRM, which is about 1-1.5%
- It will obviously result in the worse sensitivity.



# Simplifications



- To calculate that probability we need to compute 28 integrals:
  - $P_t$  and  $P_z$  of incoming partons
  - 4-momenta of 6 final partons
- To reduce CPU time, we made some assumptions:
  - $P_t$  of the incoming partons is 0. -2 integrals
  - All quarks except top are massless. -8 integrals
  - Partons and jets have the same direction. -12 integrals
  - $W$ 's and top's are on shell. -4 integrals
- Only 2 integrals in total. We'll do more in the future.



# Improvements



- Some of the improvements made:
  - Added events with 7 jets, considering last 3 jets in  $E_t$  as extra jets from radiation
    - results in better signal acceptance
  - Used refined binning for transfer functions, both in  $E_t$  in  $E_t$ 
    - results in better signal templates
- Both improvements should result in better sensitivity



# Details



Jet-parton assignments

Differential xsection

$P_t$  of ttbar system

$$P(j | M_{top}) = \frac{1}{\sigma(M_{top}) \epsilon(M_{top}) N_{combi}} \sum_{combi} \int dz_a dz_b f(z_a) f(z_b) d\sigma(M_{top}, p) TF(j | p) P_T(p)$$

Normalization factor

Integration over PDF's

Transfer functions



# dσ calculation



$$d\sigma(M_{top}, p) = \frac{|\mathcal{M}(M_{top}, p)|^2}{4E_a E_b |v_a - v_b|} (2\pi)^4 \delta^4(p_a + p_b - \sum_{i=1}^6 p_i) \prod_{i=1}^6 \frac{d^3 p_i}{(2\pi)^3 2E_i}$$

$$\mathcal{M}(M_{top}, p) \approx \frac{\bar{v}(p_{\bar{q}}) \gamma^\mu u(p_q)}{(p_q + p_{\bar{q}})^2} \cdot \frac{\bar{u}(p_u) \gamma^\alpha (1 - \gamma^5) v(p_{\bar{d}})}{P_{W^+}^2 - m_W^2 + im_W \Gamma_W} \cdot \frac{\bar{u}(p_d) \gamma^\sigma (1 - \gamma^5) v(p_{\bar{u}})}{P_{W^-}^2 - m_W^2 + im_W \Gamma_W} \cdot \bar{u}(p_b) \gamma_\alpha (1 - \gamma^5) \frac{\not{p}_t + m_t}{p_t^2 - m_t^2 + im_t \Gamma_t} \gamma_\mu \frac{\not{p}_{\bar{t}} + m_t}{p_{\bar{t}}^2 - m_t^2 + im_t \Gamma_t} \gamma_\sigma (1 - \gamma^5) v(p_{\bar{b}})$$

- We use uubar --> 6 exact tree level ME
- Spin-correlations are included
- We compute the amplitudes directly using explicit Dirac matrices and spinors

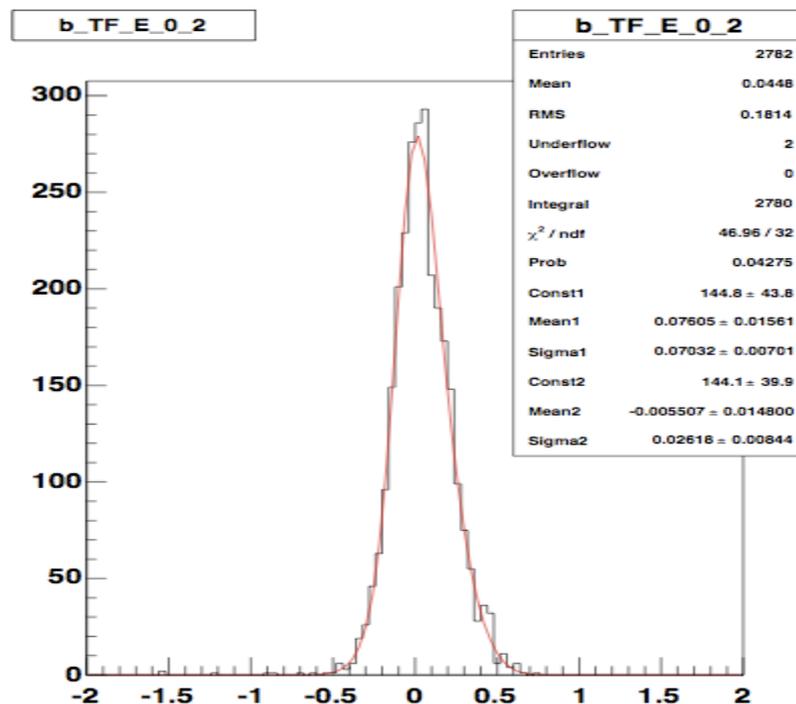


# Transfer functions



- From MC calculate the probability density function  $TF(E_j|E_p)$ 
  - $\xi = 1 - E_{jet} / E_{parton}$
- Use different TF's for different regions in  $\eta$ , energy, quark types

Example of b-quark  
Transfer Function for  
 $1.3 \leq |\eta| \leq 2$





# Prereqs effects



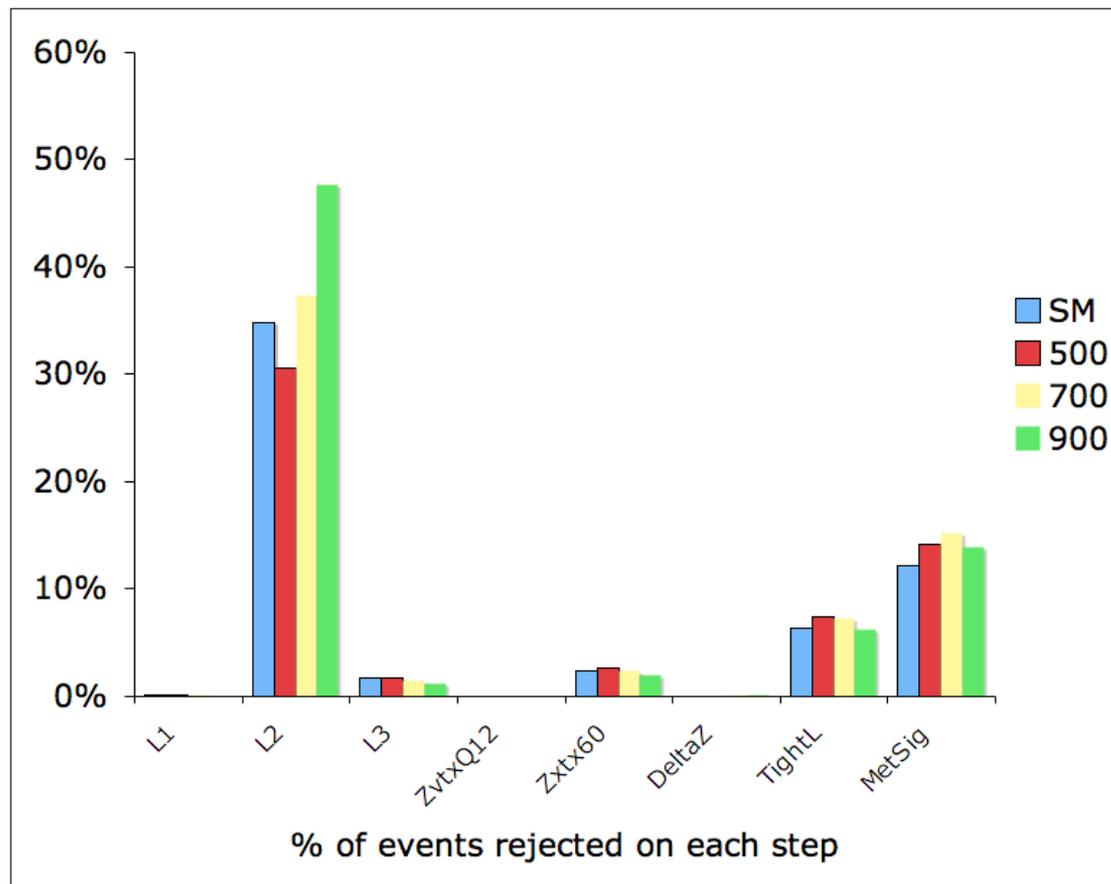
- Total efficiencies:  $\varepsilon_{tt}=42\%$ ,  $\varepsilon_{500}=43\%$ ,  $\varepsilon_{700}=36\%$ ,  $\varepsilon_{900}=28\%$
- Where do we lose events for high masses?

More interesting: Why?

Most of the events are lost on L2, which requires at least 4 clusters

For higher resonance masses, decay products are boosted more=> higher chance to merge in one cluster

See backup slides for details

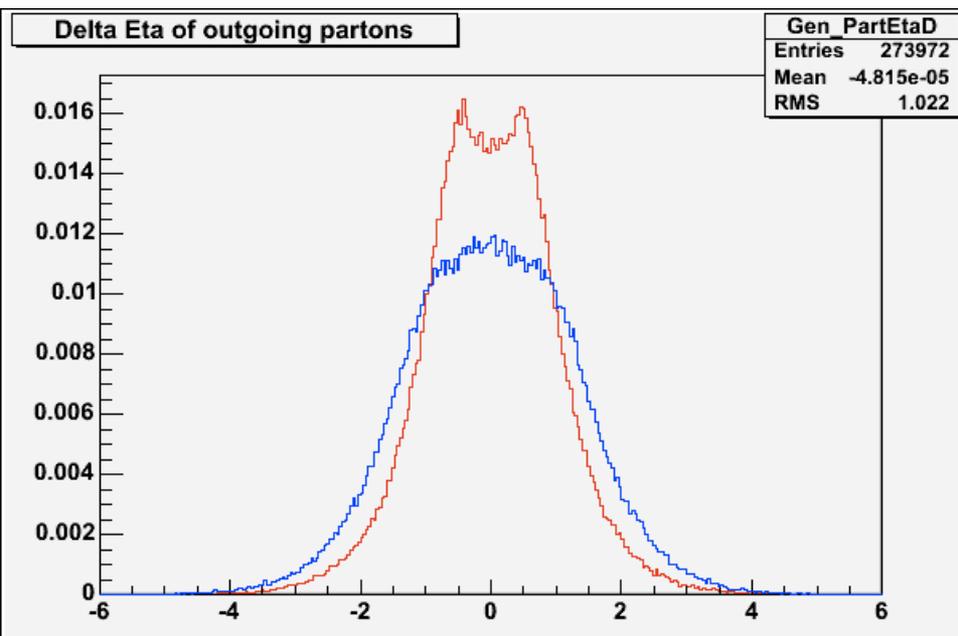
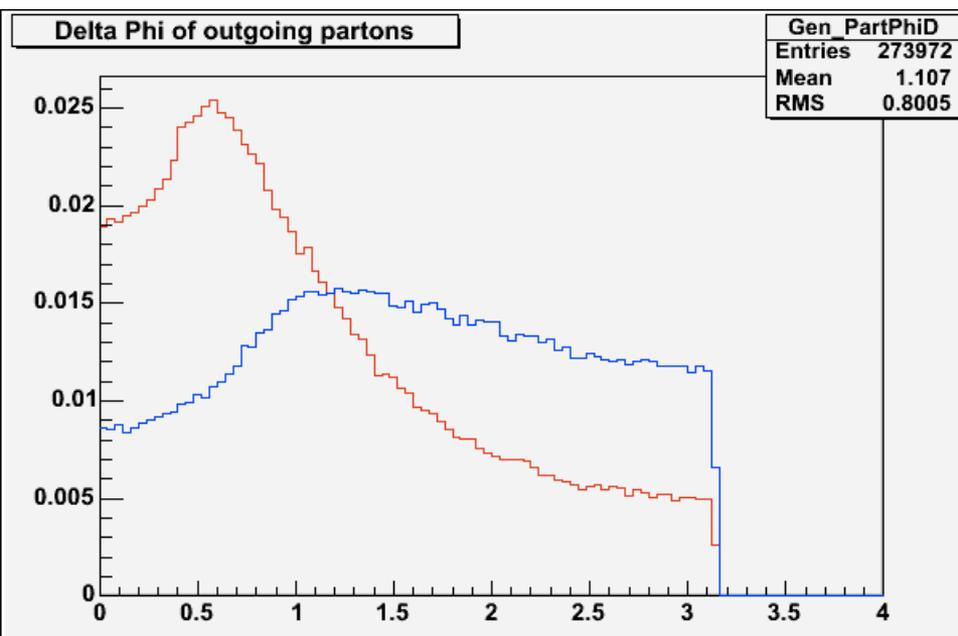




# Support for L2 issue

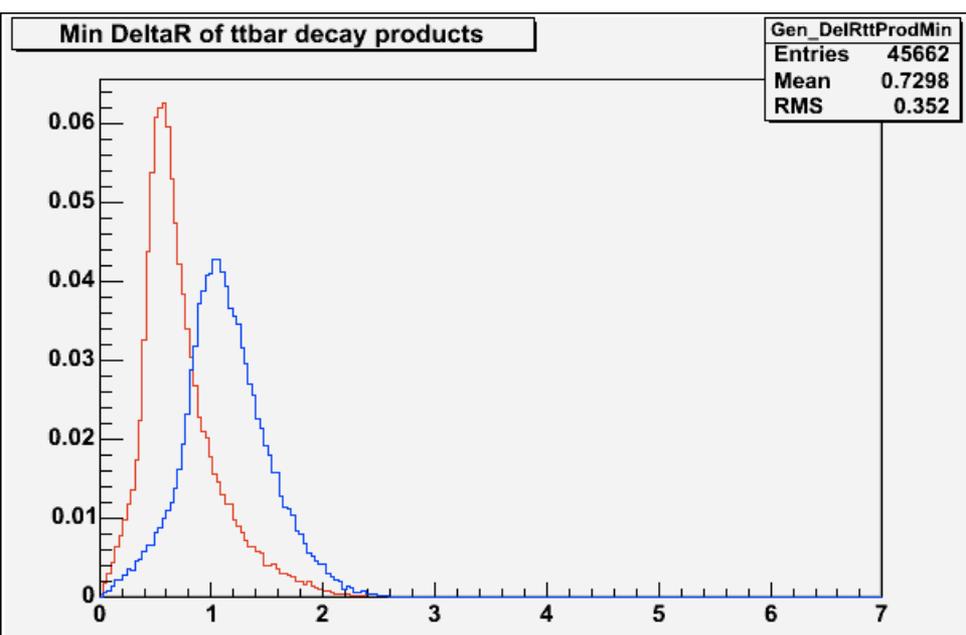
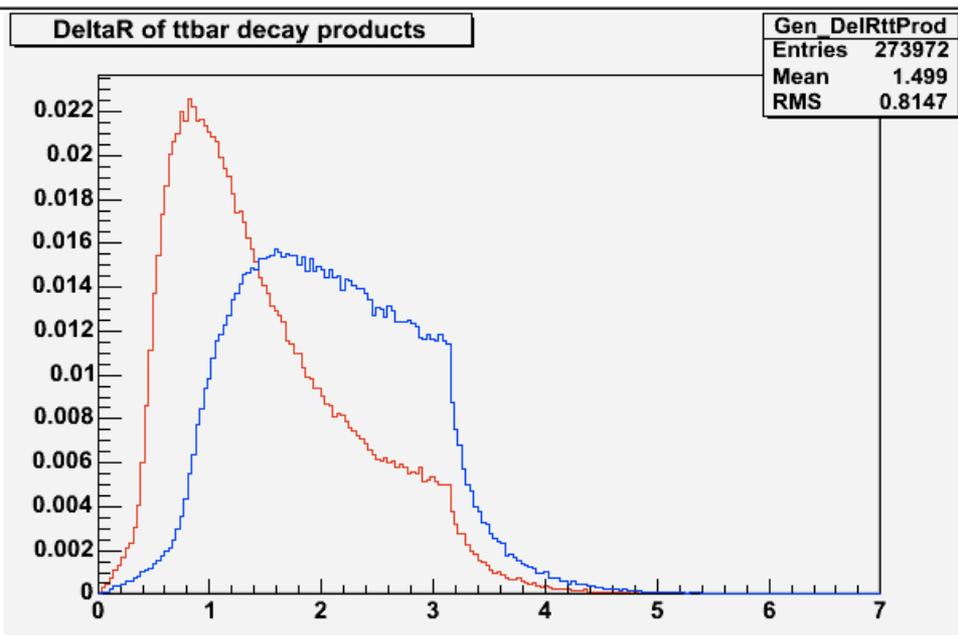


Blue 500 GeV resonance  
Red 900 GeV resonance.



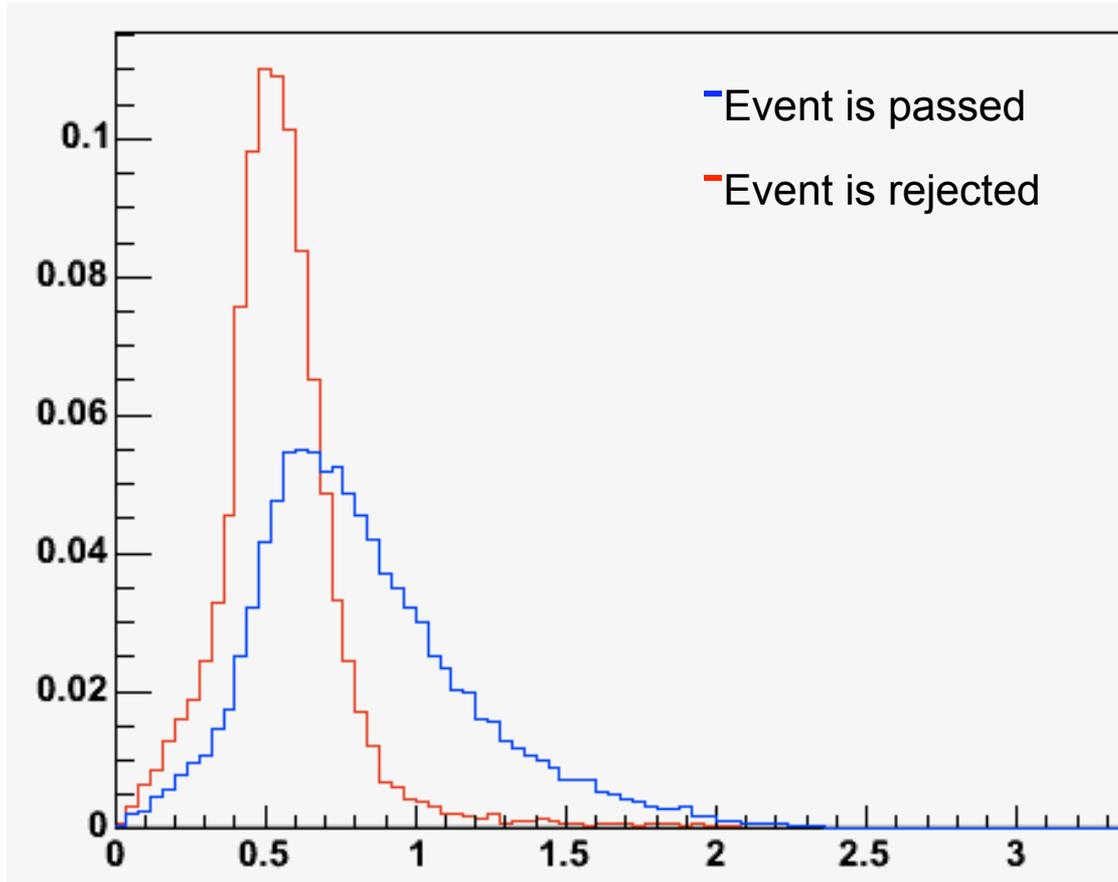


# L2 continued





# L2 final

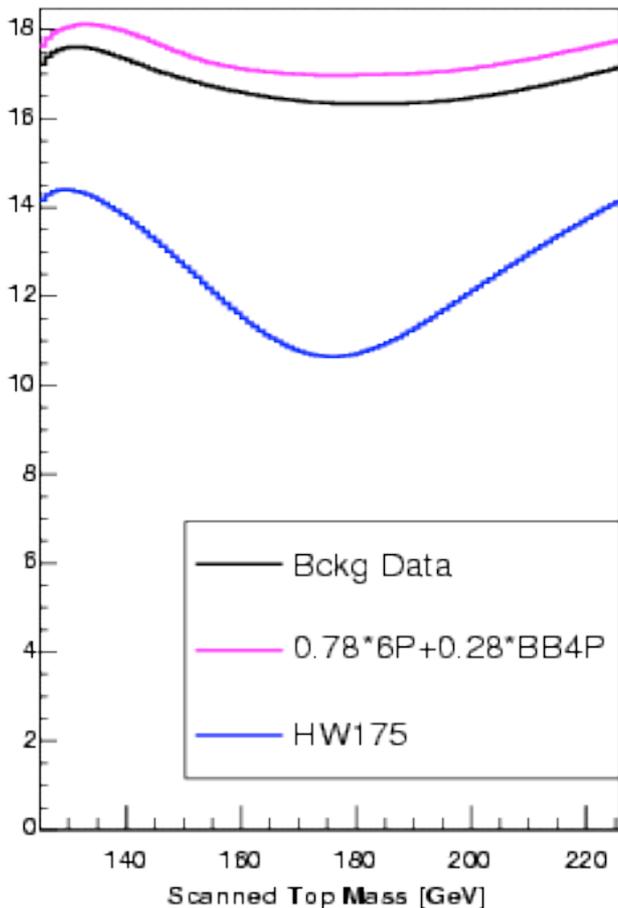




# FlaME Variable



LKLvs Mass



FlaME gives the probability of an event to come from SM  $t\bar{t}$ . Let's take advantage of it!

Here we plot  $-\log(P)$  vs top mass for various samples. As you see there is a difference between  $t\bar{t}$  and QCD

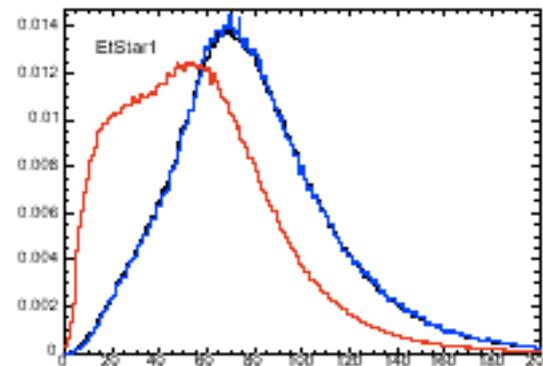
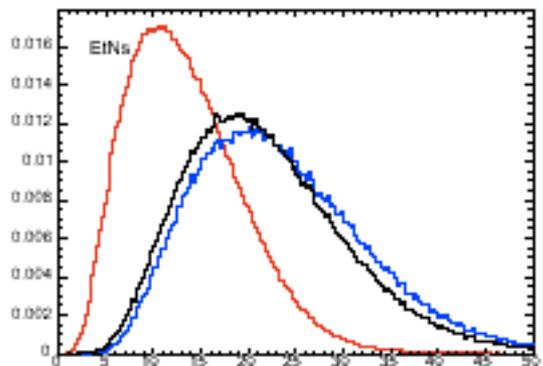
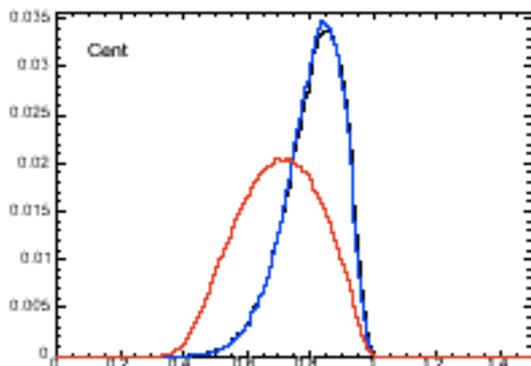
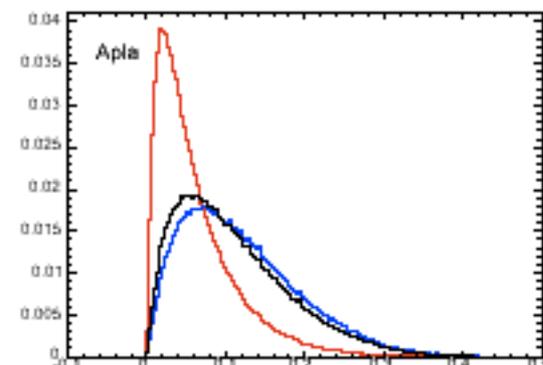
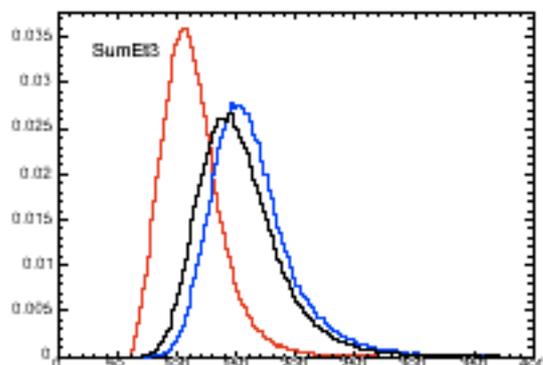
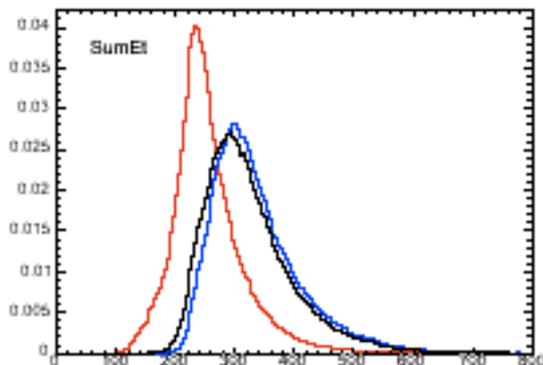
Lets calculate  $-\log(P)$  for 9 mass points: 155, 160... 195 GeV. Decided to use their sum



# ...and their distributions

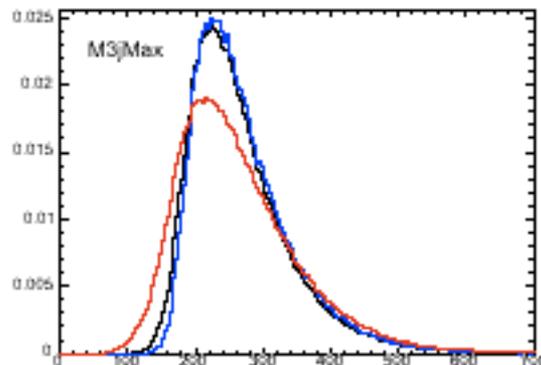
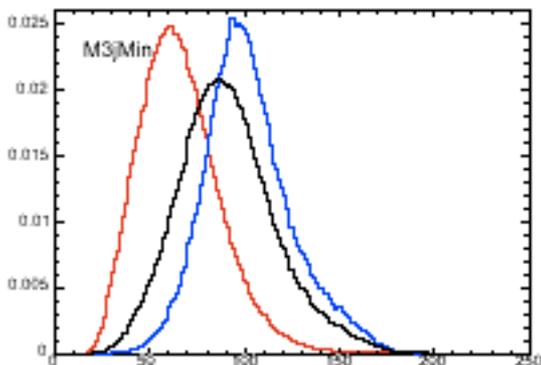
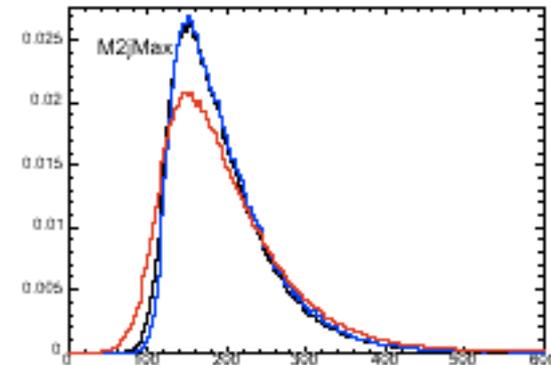
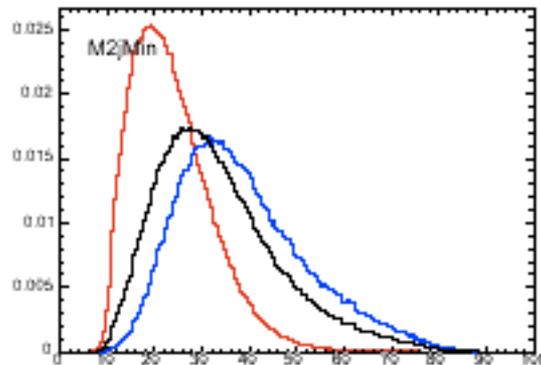
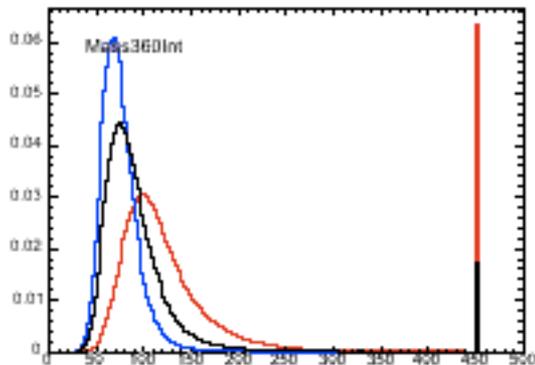


Red lines correspond to data.  
Black lines correspond to SMtt  
Blue lines correspond to SMtt matched only





# ...and their distributions





# Plug&Play



Black with FlaME, Red without FlaME, green kin. ev. sel.

